

## Anomalous $W/Z$ Quartic Couplings (QGCs)

Revised March 2024 by M.W. Grünewald (U. College Dublin) and A. Gurtu (CERN; TIFR Mumbai).

Quartic couplings,  $WWZZ$ ,  $WWZ\gamma$ ,  $WW\gamma\gamma$ , and  $ZZ\gamma\gamma$ , were studied at LEP and Tevatron at energies at which the Standard Model predicts negligible contributions to multiboson production. Thus, to parametrize limits on these couplings, an effective theory approach is adopted which supplements the Standard Model Lagrangian with higher dimensional operators which include quartic couplings. The LEP collaborations chose the lowest dimensional representation of operators (dimension 6) which presumes the  $SU(2)\times U(1)$  gauge symmetry is broken by means other than the conventional Higgs scalar doublet [1–3]. In this representation possible quartic couplings,  $a_0, a_c, a_n$ , are expressed in terms of the following dimension-6 operators [1,2];

$$\begin{aligned} L_6^0 &= -\frac{e^2}{16\Lambda^2} a_0 F^{\mu\nu} F_{\mu\nu} \vec{W}^\alpha \cdot \vec{W}_\alpha \\ L_6^c &= -\frac{e^2}{16\Lambda^2} a_c F^{\mu\alpha} F_{\mu\beta} \vec{W}^\beta \cdot \vec{W}_\alpha \\ L_6^n &= -i\frac{e^2}{16\Lambda^2} a_n \epsilon_{ijk} W_{\mu\alpha}^{(i)} W_\nu^{(j)} W^{(k)\alpha} F^{\mu\nu} \\ \tilde{L}_6^0 &= -\frac{e^2}{16\Lambda^2} \tilde{a}_0 F^{\mu\nu} \tilde{F}_{\mu\nu} \vec{W}^\alpha \cdot \vec{W}_\alpha \\ \tilde{L}_6^n &= -i\frac{e^2}{16\Lambda^2} \tilde{a}_n \epsilon_{ijk} W_{\mu\alpha}^{(i)} W_\nu^{(j)} W^{(k)\alpha} \tilde{F}^{\mu\nu} \end{aligned}$$

where  $F, W$  are photon and  $W$  fields,  $L_6^0$  and  $L_6^c$  conserve  $C, P$  separately ( $\tilde{L}_6^0$  conserves only  $C$ ) and generate anomalous  $W^+W^-\gamma\gamma$  and  $ZZ\gamma\gamma$  couplings,  $L_6^n$  violates  $CP$  ( $\tilde{L}_6^n$  violates both  $C$  and  $P$ ) and generates an anomalous  $W^+W^-Z\gamma$  coupling, and  $\Lambda$  is an energy scale for new physics. For the  $ZZ\gamma\gamma$  coupling the  $CP$ -violating term represented by  $L_6^n$  does not contribute. These couplings are assumed to be real and to vanish at tree level in the Standard Model.

Within the same framework as above, a more recent description of the quartic couplings [3] treats the anomalous parts of the  $WW\gamma\gamma$  and  $ZZ\gamma\gamma$  couplings separately, leading to two sets parametrized as  $a_0^V/\Lambda^2$  and  $a_c^V/\Lambda^2$ , where  $V = W$  or  $Z$ .

With the discovery of a Higgs at the LHC in 2012, it is then useful to go to the next higher dimensional representation (dimension 8 operators) in which the gauge symmetry is broken by the conventional Higgs scalar doublet [3,4]. There are 14 operators which can contribute to the anomalous quartic coupling signal. Some of the operators have analogues in the dimension 6 scheme. The CMS collaboration, [5], have used this parametrization, in which the connections between the two schemes are also summarized:

$$\begin{aligned} \mathcal{L}_{AQGC} &= -\frac{e^2}{8} \frac{a_0^W}{\Lambda^2} F_{\mu\nu} F^{\mu\nu} W^{+a} W_a^- \\ &\quad -\frac{e^2}{16} \frac{a_c^W}{\Lambda^2} F_{\mu\nu} F^{\mu\alpha} (W^{+\nu} W_a^- + W^{-\nu} W_a^+) \end{aligned}$$

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$$\begin{aligned}
& - e^2 g^2 \frac{\kappa_0^W}{\Lambda^2} F_{\mu\nu} Z^{\mu\nu} W^{+a} W_a^- \\
& - \frac{e^2 g^2}{2} \frac{\kappa_c^W}{\Lambda^2} F_{\mu\nu} Z^{\mu a} (W^{+\nu} W_a^- + W^{-\nu} W_a^+) \\
& + \frac{f_{T,0}}{\Lambda^4} \text{Tr}[\widehat{W}_{\mu\nu} \widehat{W}^{\mu\nu}] \times \text{Tr}[\widehat{W}_{\alpha\beta} \widehat{W}^{\alpha\beta}]
\end{aligned}$$

The energy scale of possible new physics is  $\Lambda$ , and  $g = e/\sin(\theta_W)$ ,  $e$  being the unit electric charge and  $\theta_W$  the Weinberg angle. The field tensors are described in [3,4].

The two dimension 6 operators  $a_0^W/\Lambda^2$  and  $a_c^W/\Lambda^2$  are associated with the  $WW\gamma\gamma$  vertex. Among dimension 8 operators,  $\kappa_0^W/\Lambda^2$  and  $\kappa_c^W/\Lambda^2$  are associated with the  $WWZ\gamma$  vertex, whereas the parameter  $f_{T,0}/\Lambda^4$  contributes to both vertices. There is a relationship between these two dimension 6 parameters and the dimension 8 parameters  $f_{M,i}/\Lambda^4$  as follows [3]:

$$\frac{a_0^W}{\Lambda^2} = -\frac{4M_W^2}{g^2} \frac{f_{M,0}}{\Lambda^4} - \frac{8M_W^2}{g'^2} \frac{f_{M,2}}{\Lambda^4}$$

$$\frac{a_c^W}{\Lambda^2} = -\frac{4M_W^2}{g^2} \frac{f_{M,1}}{\Lambda^4} - \frac{8M_W^2}{g'^2} \frac{f_{M,3}}{\Lambda^4}$$

where  $g' = e/\cos(\theta_W)$  and  $M_W$  is the invariant mass of the  $W$  boson. This relation provides a translation between limits on dimension 6 operators  $a_{0,c}^W$  and  $f_{M,j}/\Lambda^4$ . It is further required [4] that  $f_{M,0} = 2f_{M,2}$  and  $f_{M,1} = 2f_{M,3}$  which suppresses contributions to the  $WWZ\gamma$  vertex. The complete set of Lagrangian contributions as presented in [4] corresponds to 19 anomalous couplings in total –  $f_{S,i}$ ,  $i = 1, 2$ ,  $f_{M,i}$ ,  $i = 0, \dots, 8$  and  $f_{T,i}$ ,  $i = 0, \dots, 9$  – each scaled by  $1/\Lambda^4$ .

Another approach to couplings is the so called K-matrix framework [7], in which the anomalous couplings can be expressed in terms of two parameters  $\alpha_4$  and  $\alpha_5$ , which account for all BSM effects.

The LHC collaborations have published couplings results based on various theoretical frameworks. It is hoped that the collaborations will agree to use at least one common set of parameters to express these limits to enable the reader to make a comparison, and to allow for a possible LHC combination.

### References:

1. G. Belanger and F. Boudjema, Phys. Lett. **B288**, 201 (1992).
2. J.W. Stirling and A. Werthenbach, Eur. Phys. J. **C14**, 103 (2000);  
J.W. Stirling and A. Werthenbach, Phys. Lett. **B466**, 369 (1999);  
A. Denner *et al.*, Eur. Phys. J. **C20**, 201 (2001);  
G. Montagna *et al.*, Phys. Lett. **B515**, 197 (2001).
3. G. Belanger *et al.*, Eur. Phys. J. **C13**, 283 (2000).

4. O.J.P. Éboli, M.C. Gonzalez-Garcia, and S.M. Lietti, Phys. Rev. **D69**, 095005 (2004);  
O.J.P. Éboli, M.C. Gonzalez-Garcia, and J.K. Mizukoshi, Phys. Rev. **D77**, 073005 (2006).
5. S. Chatrchyan *et al.*, Phys. Rev. **D90**, 032008 (2014);  
S. Chatrchyan *et al.*, Phys. Rev. Lett. **114**, 051801 (2015).
6. G. Aad *et al.*, Phys. Rev. Lett. **113**, 141803 (2014).
7. A. Albateanu, W. Killian, and J. Reuter, JHEP **0811**, 010 (2008).